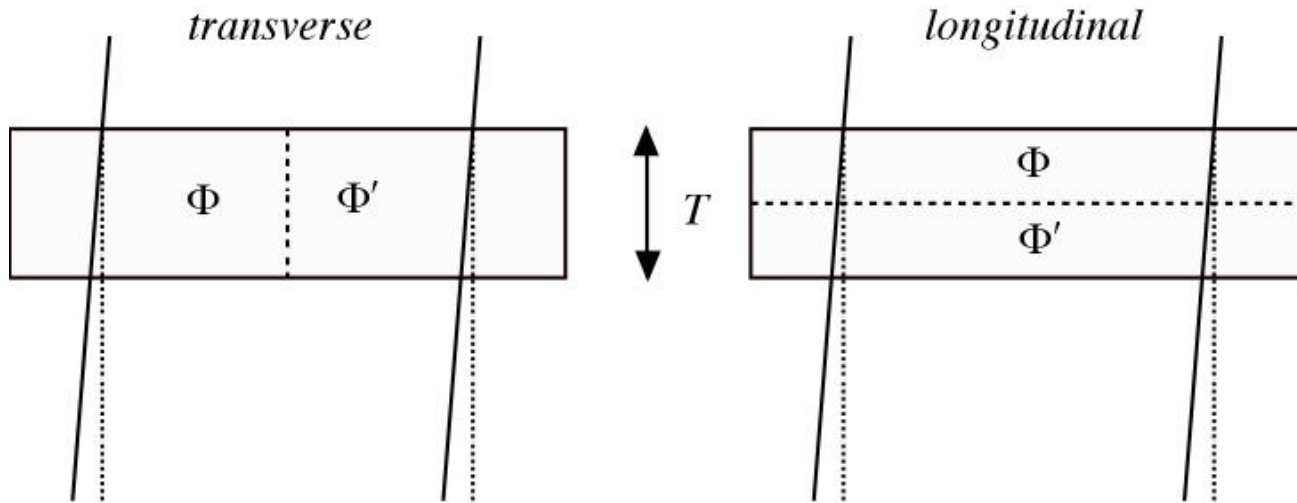


Planar defects: orientation and types



Translation:

$$\Phi'(\mathbf{r}) = \Phi(\mathbf{r} - \mathbf{R})$$

\mathbf{R} : Displacement Vector

Rotation:

$$\Phi'(\mathbf{r}) = \Phi(\tilde{M}\mathbf{r})$$

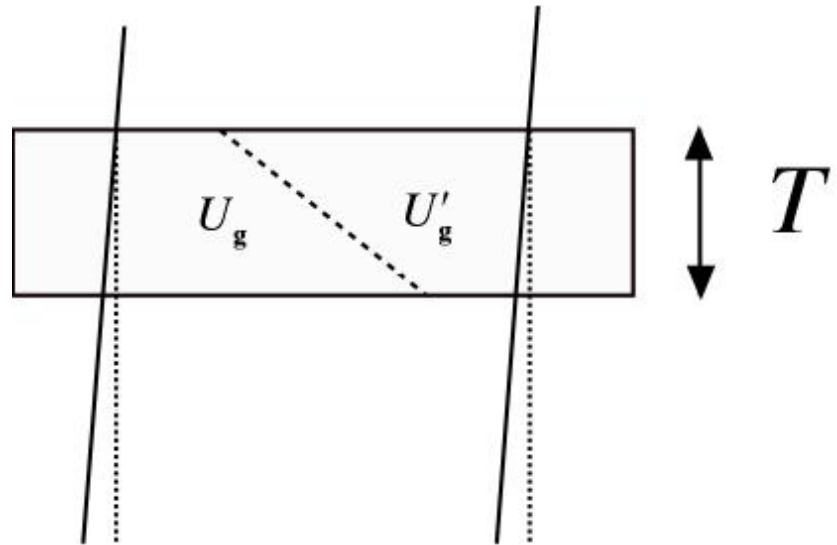
\tilde{M} : Rotation Matrix

Translation: influence on Fourier coeff's

Translation:

$$U'(\mathbf{r}) = U(\mathbf{r} - \mathbf{R})$$

$$U'(\mathbf{r}) = \sum_{\mathbf{g}} U'_g e^{2\pi i \mathbf{g} \cdot \mathbf{r}}$$



Phase factor change for each Fourier component

$$U'_g = U_g e^{-2\pi i \mathbf{g} \cdot \mathbf{R}} = U_g e^{i\alpha}$$

$$\text{APB: } U'_g = -U_g$$

$$(\alpha = 180^\circ)$$

Translation: influence on Bloch waves

Above defect:
$$\psi^{(j)}(\mathbf{r}) = e^{2\pi i \mathbf{k}^{(j)} \cdot \mathbf{r}} \cdot \sum_{\mathbf{g}} C_{\mathbf{g}}^{(j)} e^{2\pi i \mathbf{g} \cdot \mathbf{r}}$$

Below defect:
$$\begin{aligned} \psi'^{(j)}(\mathbf{r}) &= \psi^{(j)}(\mathbf{r} - \mathbf{R}) \\ &= e^{2\pi i \mathbf{k}^{(j)} \cdot (\mathbf{r} - \mathbf{R})} \cdot \sum_{\mathbf{g}} C_{\mathbf{g}}^{(j)} e^{2\pi i \mathbf{g} \cdot (\mathbf{r} - \mathbf{R})} \\ \psi'^{(j)}(\mathbf{r}) &= e^{2\pi i \mathbf{k}^{(j)} \cdot (\mathbf{r} - \mathbf{R})} \cdot \sum_{\mathbf{g}} C'_{\mathbf{g}}{}^{(j)} e^{2\pi i \mathbf{g} \cdot \mathbf{r}} \end{aligned}$$

$$C'_{\mathbf{g}}{}^{(j)} = C_{\mathbf{g}}^{(j)} e^{-2\pi i \mathbf{g} \cdot \mathbf{R}} = C_{\mathbf{g}}^{(j)} e^{i\alpha}$$

Bloch-wave coefficients are altered only if $\mathbf{g} \cdot \mathbf{R} \neq 0$

Scattering matrix (two-beam)

The dynamical diffracted intensity can be summarized using matrices:

$$\begin{pmatrix} \Psi_0(z) \\ \Psi_g(z) \end{pmatrix} = \begin{pmatrix} C_0^{(1)} & C_0^{(2)} \\ C_g^{(1)} & C_g^{(2)} \end{pmatrix} \cdot \begin{pmatrix} e^{2\pi i \gamma^{(1)} z} & 0 \\ 0 & e^{2\pi i \gamma^{(2)} z} \end{pmatrix} \cdot \begin{pmatrix} \varepsilon^{(1)} \\ \varepsilon^{(2)} \end{pmatrix} = \tilde{C} \cdot \tilde{\Gamma}(z) \cdot \begin{pmatrix} \varepsilon^{(1)} \\ \varepsilon^{(2)} \end{pmatrix}$$

$$\begin{pmatrix} \Psi_0(0) \\ \Psi_g(0) \end{pmatrix} = \tilde{C} \cdot \begin{pmatrix} \varepsilon^{(1)} \\ \varepsilon^{(2)} \end{pmatrix} \qquad \begin{pmatrix} \varepsilon^{(1)} \\ \varepsilon^{(2)} \end{pmatrix} = (\tilde{C})^{-1} \cdot \begin{pmatrix} \Psi_0(0) \\ \Psi_g(0) \end{pmatrix}$$

$$\begin{pmatrix} \Psi_0(z) \\ \Psi_g(z) \end{pmatrix} = \tilde{C} \cdot \tilde{\Gamma}(z) \cdot (\tilde{C})^{-1} \cdot \begin{pmatrix} \Psi_0(0) \\ \Psi_g(0) \end{pmatrix}$$

Scattering Matrix: $\tilde{P}(z) = \tilde{C} \cdot \tilde{\Gamma}(z) \cdot \tilde{C}^{-1}$

$$\begin{pmatrix} \Psi_0(z) \\ \Psi_g(z) \end{pmatrix} = \tilde{P}(z) \cdot \begin{pmatrix} \Psi_0(0) \\ \Psi_g(0) \end{pmatrix}$$

Propagation across a planar defect (I)

Below the fault:

$$\begin{pmatrix} \Psi'_0(z) \\ \Psi'_g(z) \end{pmatrix} = \tilde{C}' \cdot \tilde{\Gamma}'(z-t) \cdot \begin{pmatrix} \varepsilon'^{(1)} \\ \varepsilon'^{(2)} \end{pmatrix}$$

Same eigenvalues:

Different Bloch wave coefficients:

$$\Gamma'(z-t) = \Gamma(z-t) = \begin{pmatrix} e^{2\pi i \gamma^{(1)}(z-t)} & 0 \\ 0 & e^{2\pi i \gamma^{(2)}(z-t)} \end{pmatrix} \quad \tilde{C}' = \begin{pmatrix} 1 & 0 \\ 0 & e^{-i\alpha} \end{pmatrix} \cdot \tilde{C}$$

$$\begin{pmatrix} \Psi'_0(t) \\ \Psi'_g(t) \end{pmatrix} = \tilde{C}' \cdot \begin{pmatrix} \varepsilon'^{(1)} \\ \varepsilon'^{(2)} \end{pmatrix} \quad \begin{pmatrix} \varepsilon'^{(1)} \\ \varepsilon'^{(2)} \end{pmatrix} = (\tilde{C}')^{-1} \cdot \begin{pmatrix} \Psi'_0(t) \\ \Psi'_g(t) \end{pmatrix}$$

Scattering Matrix: $\tilde{P}'(z-t) = \tilde{C}' \cdot \tilde{\Gamma}'(z-t) \cdot (\tilde{C}')^{-1}$

$$\begin{pmatrix} \Psi'_0(z) \\ \Psi'_g(z) \end{pmatrix} = \tilde{P}'(z-t) \cdot \begin{pmatrix} \Psi'_0(t) \\ \Psi'_g(t) \end{pmatrix}$$

Propagation across a planar defect (II)

Boundary condition at the buried phase boundary interface ($z = t$):

$$\begin{pmatrix} \Psi'_0(t) \\ \Psi'_g(t) \end{pmatrix} = \begin{pmatrix} \Psi_0(t) \\ \Psi_g(t) \end{pmatrix} = \tilde{P}(t) \cdot \begin{pmatrix} \Psi_0(0) \\ \Psi_g(0) \end{pmatrix}$$

$$\begin{pmatrix} \Psi'_0(z) \\ \Psi'_g(z) \end{pmatrix} = \tilde{P}'(z-t) \cdot \begin{pmatrix} \Psi'_0(t) \\ \Psi'_g(t) \end{pmatrix} = \tilde{P}'(z-t) \cdot \tilde{P}(t) \cdot \begin{pmatrix} \Psi_0(0) \\ \Psi_g(0) \end{pmatrix}$$

$$\begin{pmatrix} \Psi_0(0) \\ \Psi_g(0) \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$\begin{pmatrix} \Psi'_0(T) \\ \Psi'_g(T) \end{pmatrix} = \tilde{P}'(T-t) \cdot \tilde{P}(t) \cdot \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

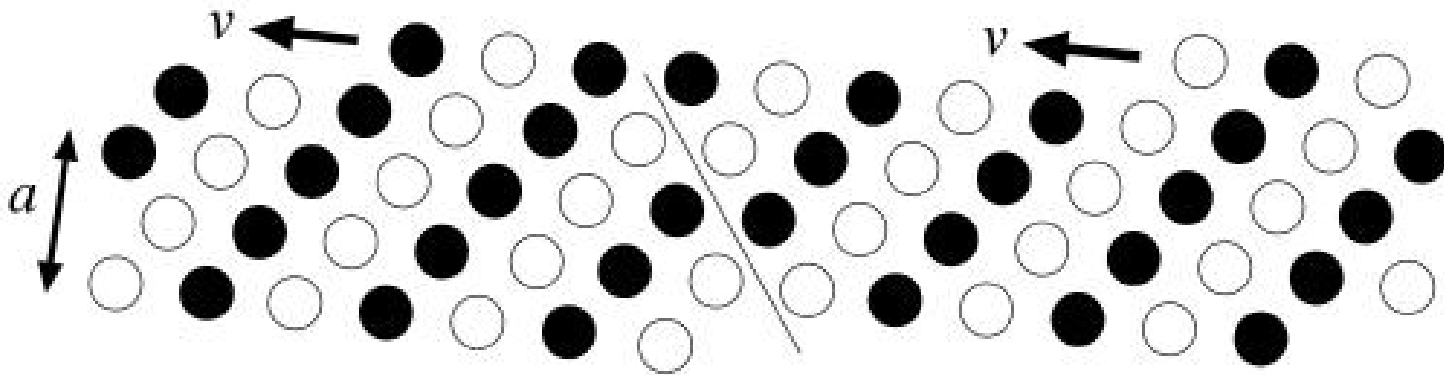
Antiphase boundary

180° phase shift in ordering sequence

$$\alpha = 180^\circ$$

$$F_g = f_A - f_B$$

$$F'_g = f_B - f_A = -F_g$$



Depiction of APB propagation via step-flow-driven crystal growth

Propagation across an APB (I)

General two-beam result:

$$\Gamma(z) = e^{\pi i w z / \xi} \cdot \begin{pmatrix} e^{\pi i \sqrt{1+w^2} z / \xi} & 0 \\ 0 & e^{-\pi i \sqrt{1+w^2} z / \xi} \end{pmatrix}$$

$$w = 1/s\xi$$

$$\sin(\beta) = 1/\sqrt{1+w^2}$$

$$\cos(\beta) = w/\sqrt{1+w^2}$$

$$\tilde{C} = \begin{pmatrix} \sin(\beta/2) & -\cos(\beta/2) \\ \cos(\beta/2) & \sin(\beta/2) \end{pmatrix} \quad \tilde{C}' = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \cdot \tilde{C} = \begin{pmatrix} \sin(\beta/2) & -\cos(\beta/2) \\ -\cos(\beta/2) & -\sin(\beta/2) \end{pmatrix}$$

Scattering matrices:

$$\tilde{P}(z) = \tilde{C} \cdot \tilde{\Gamma}(z) \cdot \tilde{C}^{-1}$$

$$\tilde{P}'(z-t) = \tilde{C}' \cdot \tilde{\Gamma}'(z-t) \cdot (\tilde{C}')^{-1}$$

Propagation across an APB (II)

Strong-beam result: $\gamma^{(1,2)} = \pm 1/2\xi$

$$\Gamma(z) = \begin{pmatrix} e^{\pi iz/\xi} & 0 \\ 0 & e^{-\pi iz/\xi} \end{pmatrix} \quad \tilde{C} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & -1 \\ 1 & 1 \end{pmatrix} \quad \tilde{C}' = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & -1 \\ -1 & -1 \end{pmatrix}$$

Above the APB: $\tilde{\Gamma}(t) = \begin{pmatrix} e^{ia} & 0 \\ 0 & e^{-ia} \end{pmatrix}$ Define: $a = \pi t/\xi$ $b = \pi T/\xi$

$$\tilde{P}(t) = \frac{1}{2} \begin{pmatrix} 1 & -1 \\ 1 & 1 \end{pmatrix} \cdot \begin{pmatrix} e^{ia} & 0 \\ 0 & e^{-ia} \end{pmatrix} \cdot \begin{pmatrix} 1 & 1 \\ -1 & 1 \end{pmatrix} = \begin{pmatrix} \cos(a) & i \sin(a) \\ i \sin(a) & \cos(a) \end{pmatrix}$$

Below the APB: $\tilde{\Gamma}(T-t) = \begin{pmatrix} e^{i(b-a)} & 0 \\ 0 & e^{-i(b-a)} \end{pmatrix}$

$$\tilde{P}'(T-t) = \frac{1}{2} \begin{pmatrix} 1 & -1 \\ -1 & -1 \end{pmatrix} \cdot \begin{pmatrix} e^{i(b-a)} & 0 \\ 0 & e^{-i(b-a)} \end{pmatrix} \cdot \begin{pmatrix} 1 & -1 \\ -1 & -1 \end{pmatrix} = \begin{pmatrix} \cos(b-a) & -i \sin(b-a) \\ -i \sin(b-a) & \cos(b-a) \end{pmatrix}$$

Propagation across an APB (III)

Apply boundary condition at entrance surface:

$$\tilde{\mathbf{P}}(t) \cdot \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} \cos(a) \\ i \sin(a) \end{pmatrix}$$

$$\tilde{\mathbf{P}}(T-t) \cdot \tilde{\mathbf{P}}(t) \cdot \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} \cos(b-a) & -i \sin(b-a) \\ -i \sin(b-a) & \cos(b-a) \end{pmatrix} \cdot \begin{pmatrix} \cos(a) \\ i \sin(a) \end{pmatrix} = \begin{pmatrix} \cos(2a-b) \\ i \sin(2a-b) \end{pmatrix}$$

Diffracted-beam amplitudes:

$$\begin{pmatrix} \Psi'_0(T) \\ \Psi'_g(T) \end{pmatrix} = \begin{pmatrix} \cos[\pi(2t-T)/\xi] \\ i \sin[\pi(2t-T)/\xi] \end{pmatrix}$$

Diffracted Intensity Across an APB

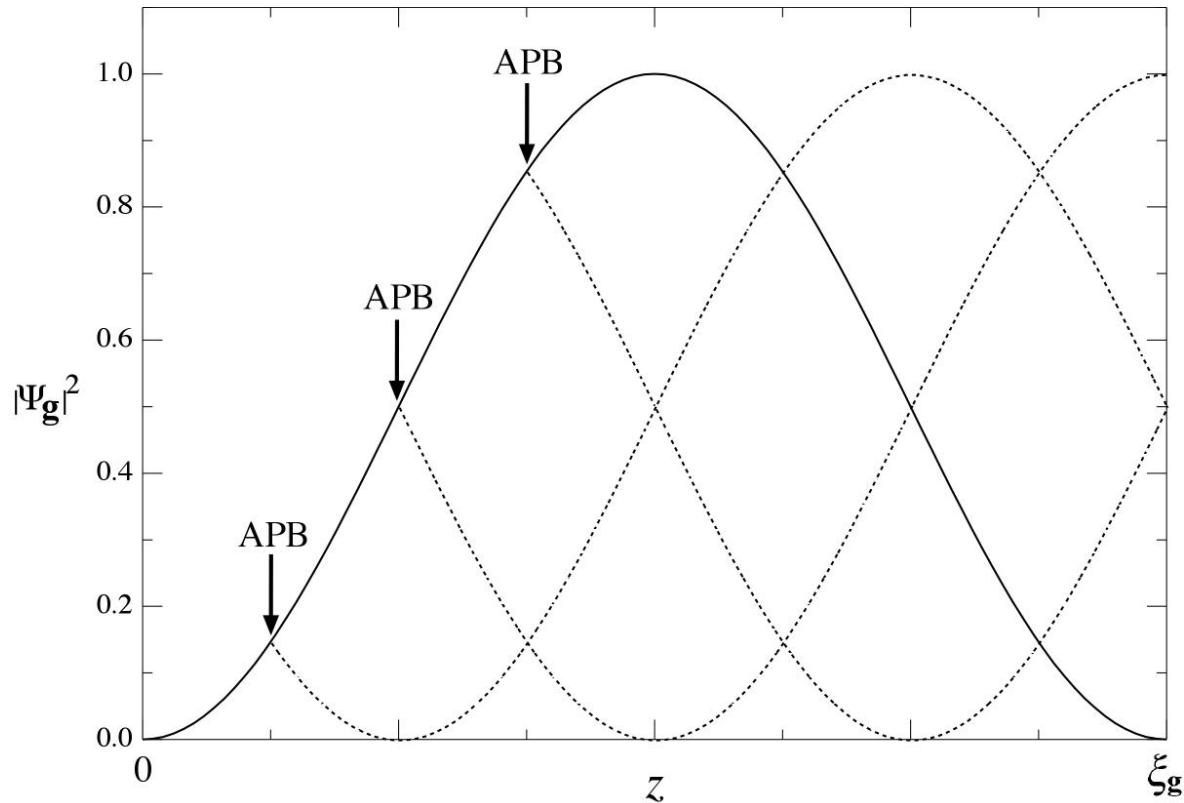
Strong-Beam Case:

Above:

$$|\Psi_{\mathbf{g}}(z)|^2 = \sin^2(\pi z/\xi)$$

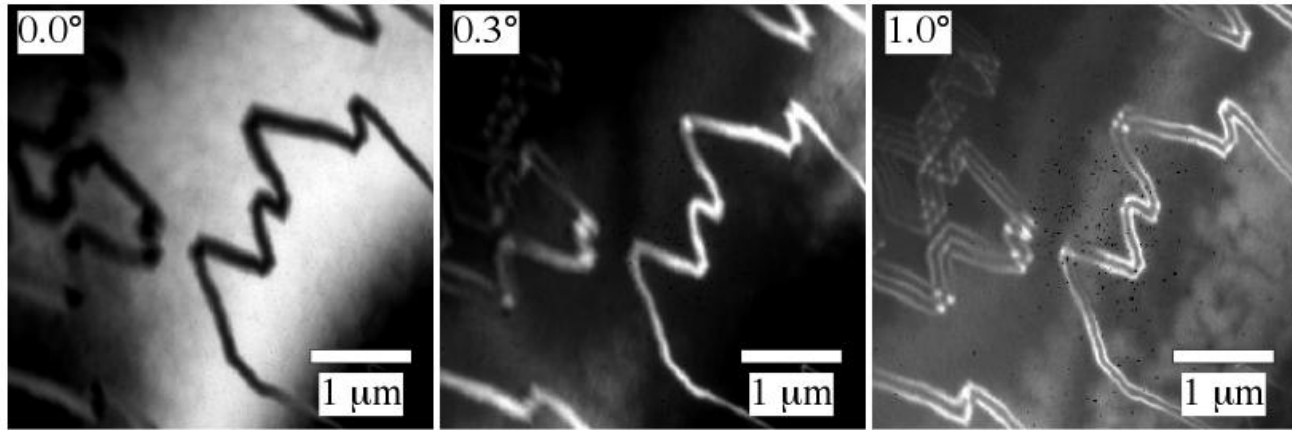
Below:

$$|\Psi'_{\mathbf{g}}(z)|^2 = \sin^2[\pi(2t - z)/\xi]$$

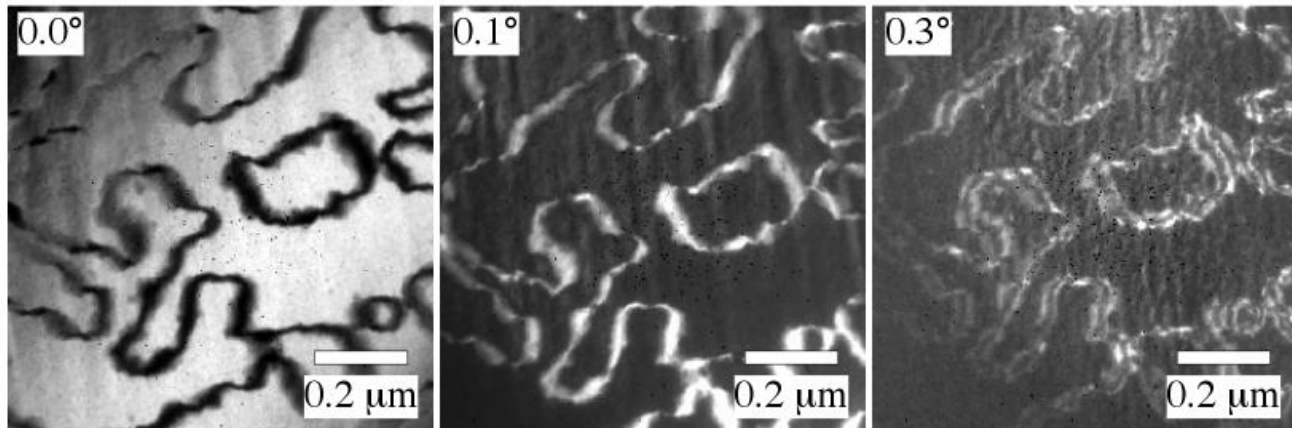


DF images of inclined APBs: influence of tilt

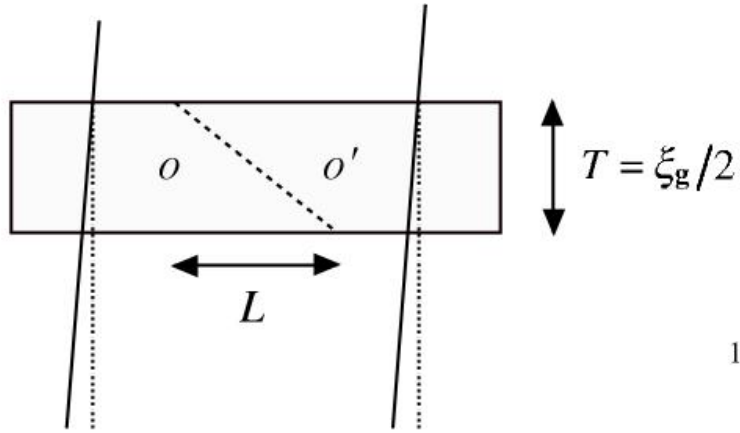
GaInP



GaInAs

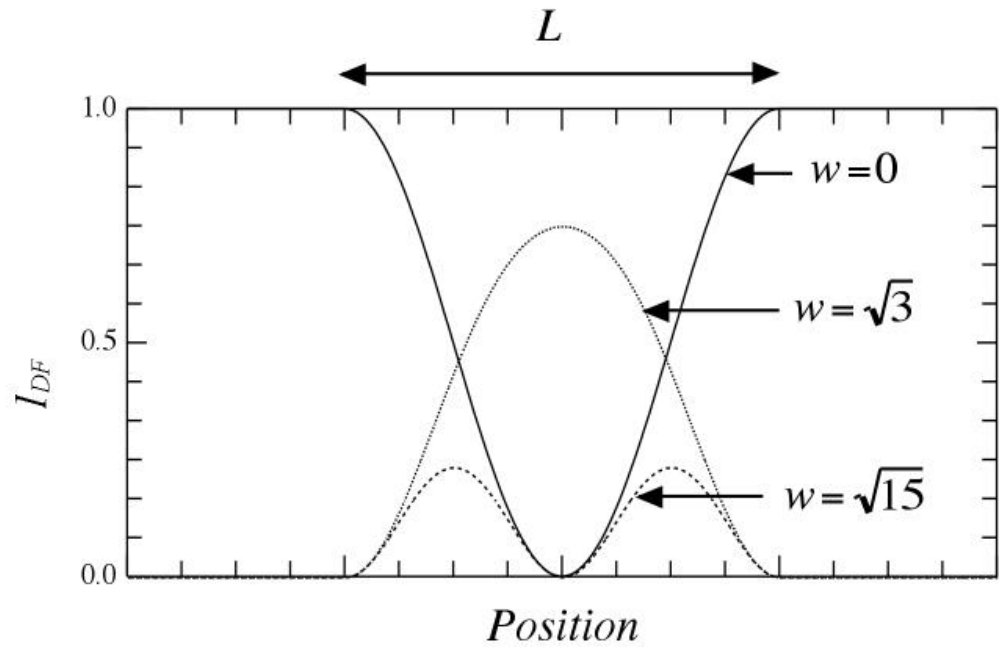
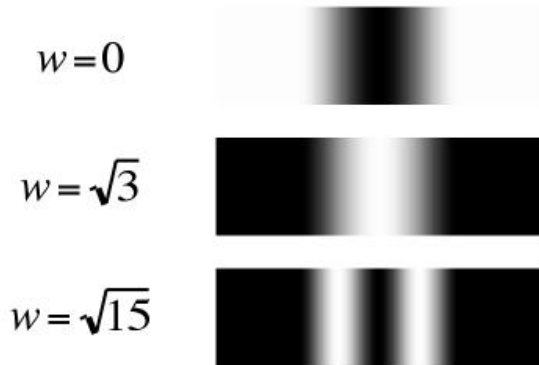


Inclined APBs: two-beam analysis



$$s_{eff} \equiv \left(\frac{1}{\xi} \right) \cdot \sqrt{w^2 + 1}$$

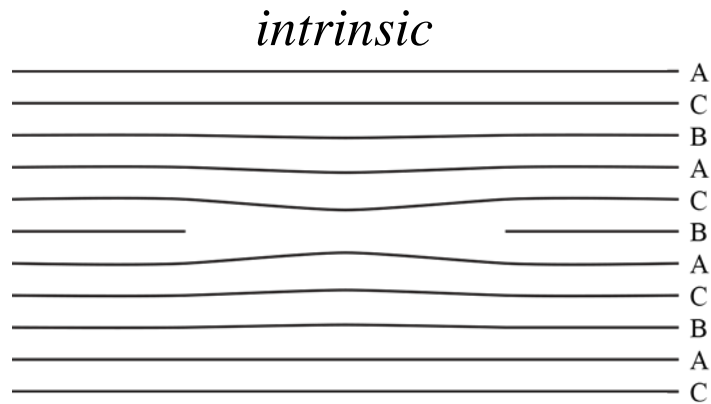
$$\omega^2 + 1 = N^2$$



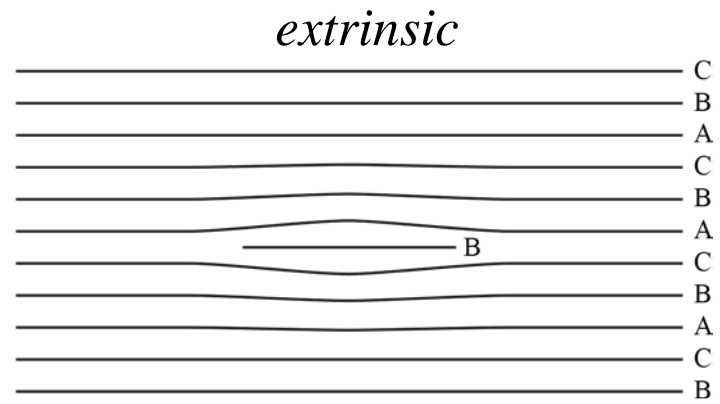
Stacking faults in fcc crystals

Two types:

Intrinsic and Extrinsic

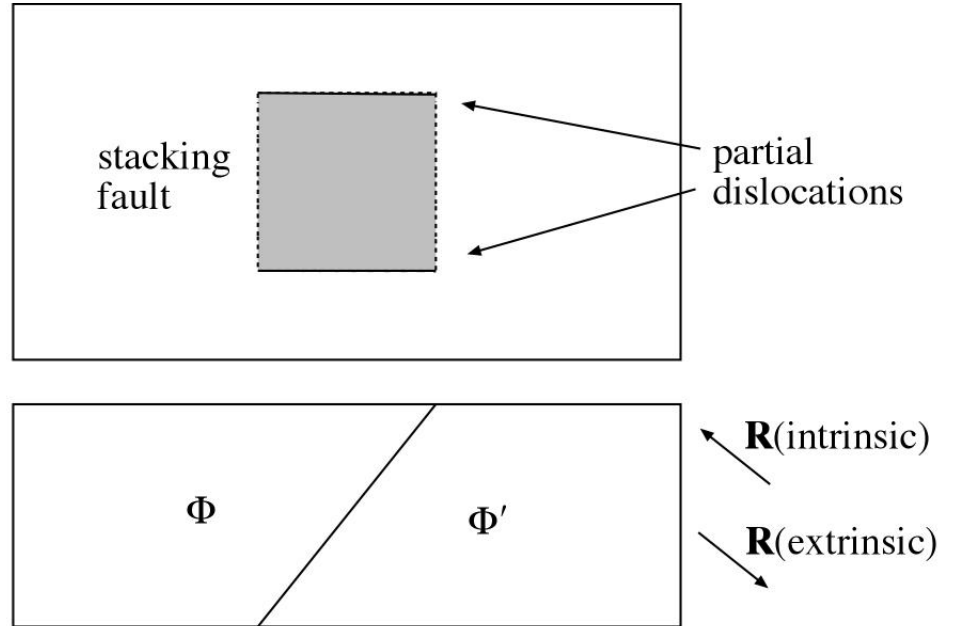


↑
[111]

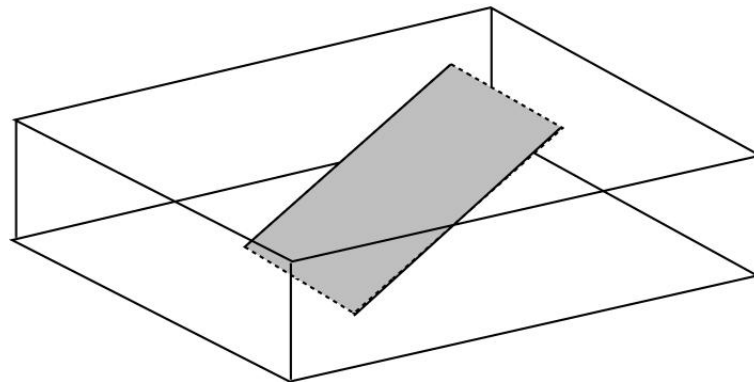


Stacking fault geometry

Partial dislocations border fault

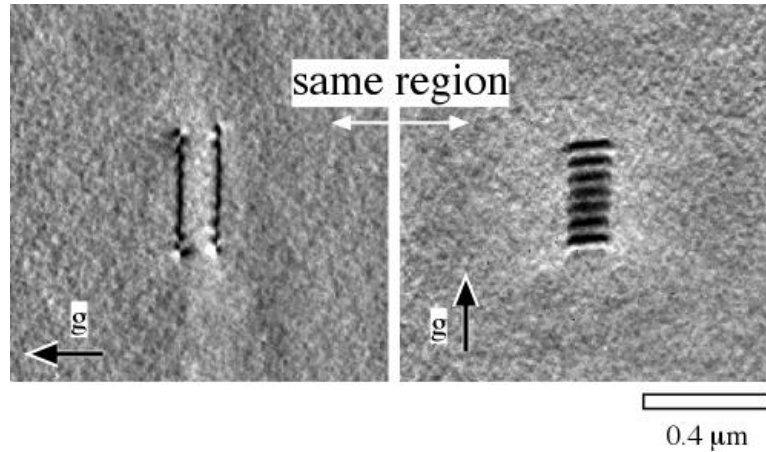


Fault is observed with $\mathbf{g} \cdot \mathbf{R} \neq 0$

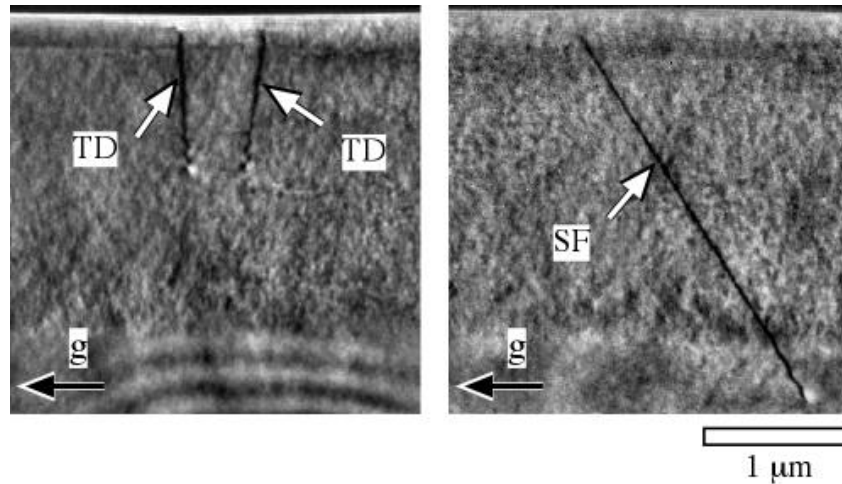


Stacking fault images

[001]
Plan view
{220} DF

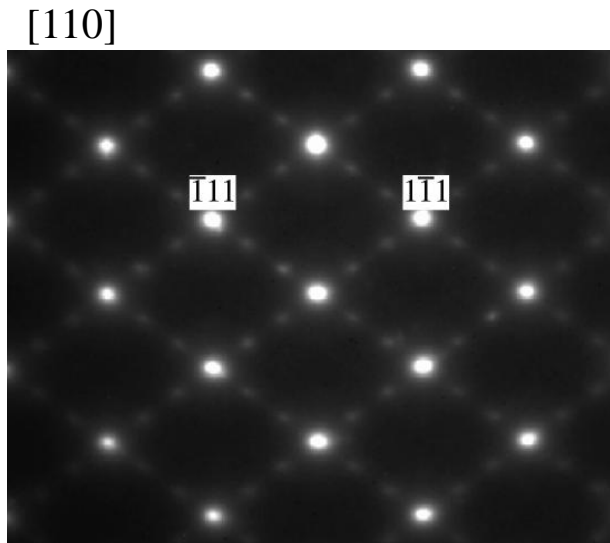


[110]
Cross section
{220} DF

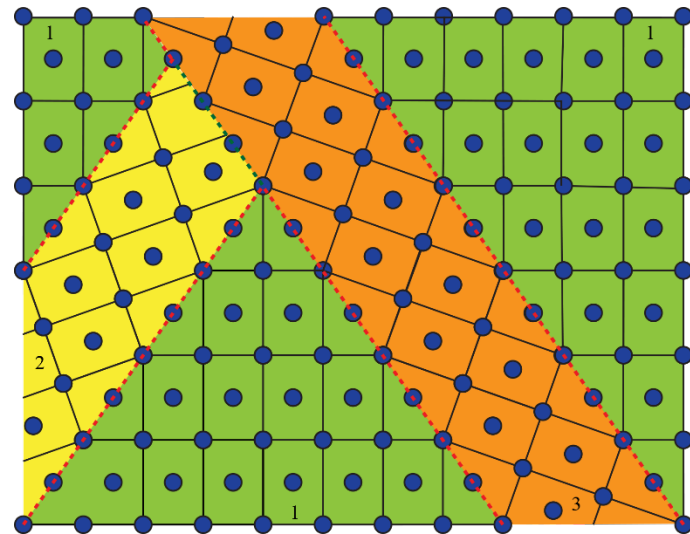
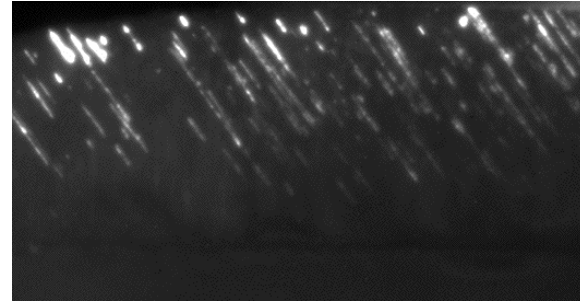


Sometimes hard to distinguish threading dislocations from partial dislocations

Twin boundaries in fcc crystals



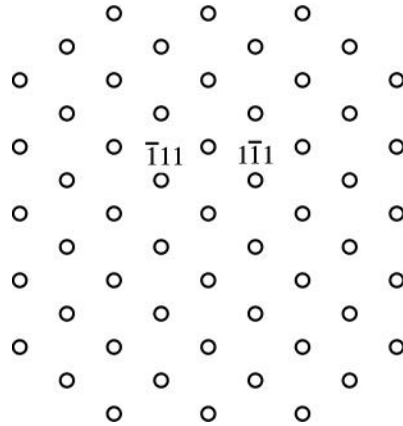
$1/3(111)$ dark-field image



Give rise to extra spots at $1/3(111)$ positions

Origin of 1/3 position spots

Matrix

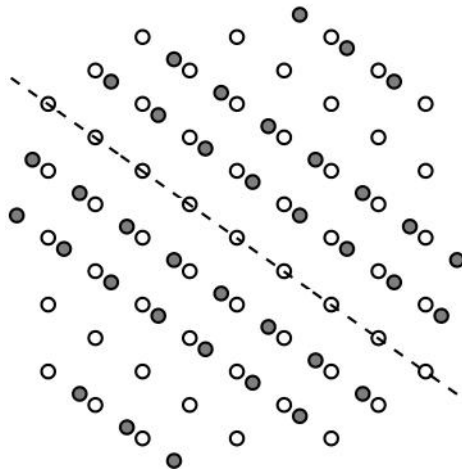


[110] fcc

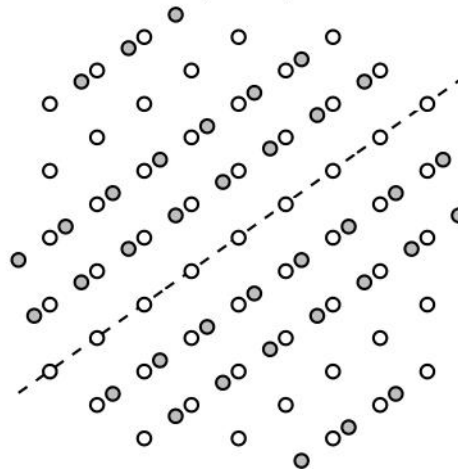
180° Rotation about (111) directions

Dynamical diffraction gives complete pattern

Matrix + ($\bar{1}11$) variant



Matrix + ($1\bar{1}1$) variant



Matrix + Both Variants

