Lecture X: The Dirac Equation and Graphene

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Introduction: In light of Einstien's theory of relativity, energy and mass of a particle are related. The energy of a particle at rest (in motion) are given as:

$$E_o = m_o c^2$$
 Einstien's equation for particle at rest (1)

$$E = \sqrt{(m_o c^2)^2 + (cp)^2}$$
 Einstien's equation for relativistic particle (2)

we can calculate the velocity v as:

$$v = \frac{\partial E}{\partial p} = c \frac{cp}{\sqrt{(m_o c^2)^2 + (cp)^2}} \tag{3}$$

From this expression, we can see the velocity v of massive particles as a function of momentum p approaches, but cannot exceed, the speed of light c:



Figure 1: Velocity v of a massive particle as a function of momentum \mathbf{p} according to equation 3.

Limiting cases of equation 2 for $m_o \to 0$ and $v \ll c$ yield:

$$E = c |\mathbf{p}|$$
 Massless particle (4)
 $E \approx m_o c^2 + p^2/2m$ Non-relativistic, massive particle (5)

which, as seen in figure 2, yield linear and parabolic dispersion relations for E vs. **p**:



Figure 2: Energy E vs. momentum **p** for (left): a massless particle and (right): a massive particle according to equations 4-5.

To reconcile the relativistic energy equation 2 with quantum mechanics, we need to go back to the operator correspondence:

$$E \longrightarrow i\hbar \frac{\partial}{\partial t}$$
 and $\mathbf{p} \longrightarrow -i\hbar \nabla$

from which, we recall, we get Schrödingers equation for a free particle immediately from the definition of the Hamiltonian (V = 0 here):

$$i\hbar\frac{\partial\Psi}{\partial t}=-\frac{\hbar^2}{2m}\nabla^2\Psi$$

For a relativistic particle, we need to take into account the rest mass. Following a similar recipe but using the expression for the energy of a relativistic free particle given by equation 2, we have (here, $m \equiv m_o$):

$$E^{2} = c^{2}p^{2} + m^{2}c^{4} \longrightarrow -\hbar^{2}\frac{\partial^{2}\Psi}{\partial t^{2}} = \left(-\hbar^{2}c^{2}\nabla^{2} + m^{2}c^{4}\right)\Psi$$
(6)

This is the Klein Gordon equation. However, we need a form which is first order in time. To this end, Dirac found the kinetic energy terms can be factored, and that the proper representation is based on a two dimensional basis, which treats spin up and spin down particles separately:

$$(p_x^2 + p_y^2 + m^2) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = \begin{pmatrix} m & p_x - ip_y \\ p_x + ip_y & -m \end{pmatrix}^2$$

Taking the square root of this expression gives a linear form for the energy operator, we can now use the same correspondence to write a first order quantum mechanical equation of motion for a relativistic particle:

$$i\hbar\frac{\partial\Psi}{\partial t} = \left[-i\hbar\left(\sigma_x\frac{\partial}{\partial x} + \sigma_y\frac{\partial}{\partial y}\right) + \sigma_z m\right]\Psi\tag{7}$$

which is the Dirac equation, where:

$$\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \text{ and } \Psi = \begin{pmatrix} \Psi_{\uparrow} \\ \Psi_{\downarrow} \end{pmatrix}$$

The Dirac equation has several significant consequences, for instance, the existence of anti-particles and spin. As seen in the dispersion relation for graphene, for low energies near the Dirac point, electrons obey a Dirac equation with m = 0 and $c = v_F$, the Fermi velocity. We say the charge carriers in this case are "emergent" Dirac Fermions, with the two subtypes A(B) corresponding to two spin spin states in the Dirac equation (termed "psuedo-spin"):

$$\boldsymbol{\Psi} = \begin{pmatrix} \Psi_{\uparrow} \\ \Psi_{\downarrow} \end{pmatrix} = \begin{pmatrix} \Psi_A \\ \Psi_B \end{pmatrix}$$

Because of these emergent properties, charge carriers in graphene are expected to have unique properties predicted by the theory, *e.g.* extraordinary transport properties, which have been borne out by experiments.